

Construction and Analysis of Consistent Atomistic/Continuum Coupling Methods

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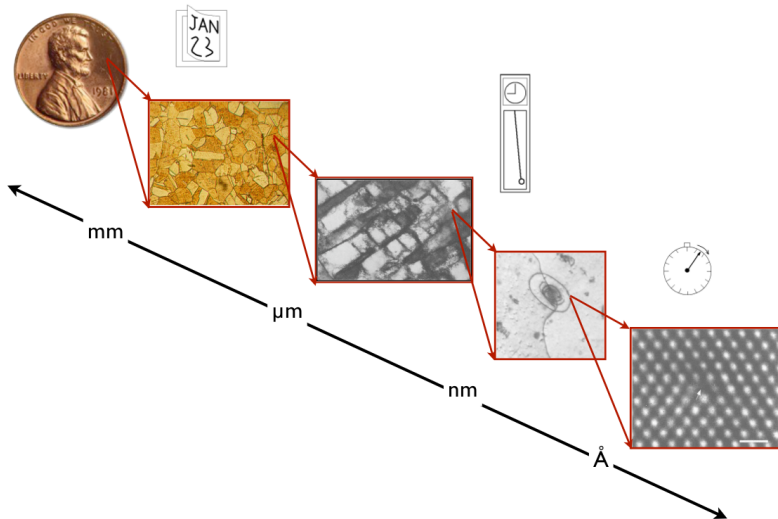
Dundee, Sep 9, 2013

Outline

In this talk, we will focus on multiscale methods for 2D/3D point defects at 0T, namely, the coupling of length scales.

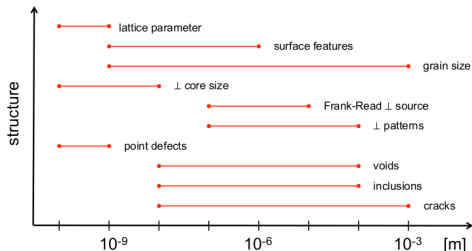
- Introduction
 - Multiscales in Materials & Crystalline Defects
 - Atomistic Simulation & Its Continuum Coarse Graining
 - Atomistic/Continuum Coupling
 - Issue of Patch Test Consistent (Ghost Force)
- Construction of Consistent Method
 - Overview
 - First Order Consistency
 - Construction
 - Numerical Experiment
 - Optimized construction
- Stability
 - Overview
 - Universally Stable Method
 - Stability Gap and Stabilization
 - Balance between Consistency and Stability
- Outlook

Multiple Scales in Materials



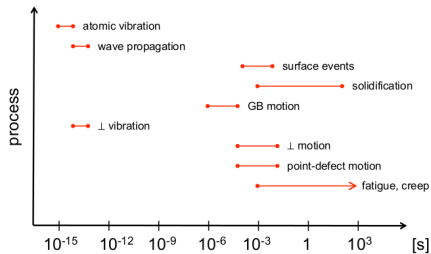
[Modeling Materials, Tadmor & Miller 2012]

Multiple Scales in Materials



Structural features:

- range in size from 10^{-10} to 10^{-3} m
- interact atomistically at short distances and over long distances via long-range elastic stress fields.



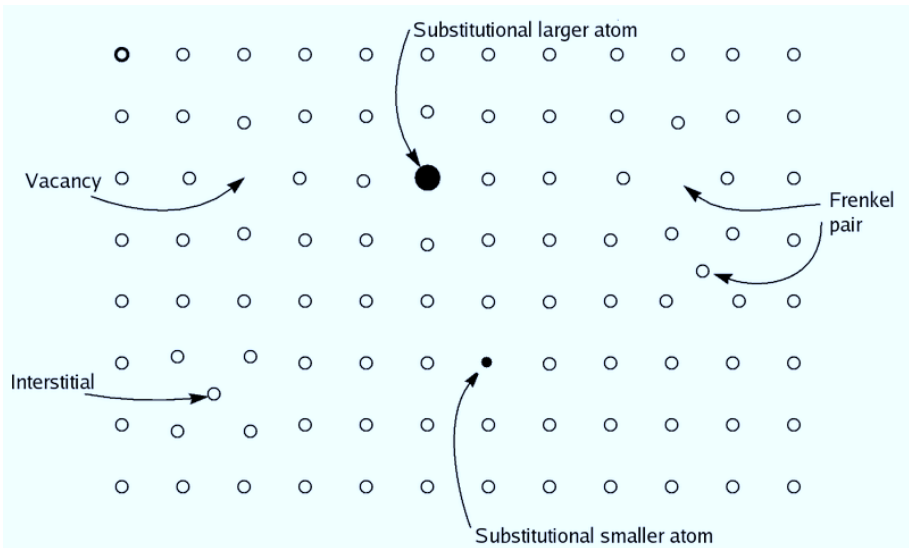
Materials processes:

- occur over time scales ranging from 10^{-15} to years.
- this broad range can sometimes be an asset when scale separation occurs.

Modeling material behavior requires methods able to span across length and time scales.

[Modeling Materials, Tadmor & Miller 2012]

Point Defects in 2D



[Wikipedia]

Atomistic Mechanics (0T statics)

- **Atomistic body:** N atoms at positions $y = (y_n)_{n=1}^N \in \mathbb{R}^{d \times N}$
- **Total energy** of configuration y :

$$\min \mathcal{E}_a^{\text{tot}}(y) := \mathcal{E}^a(y) + \mathcal{P}_a(y)$$

- \mathcal{E}^a = interaction energy, \mathcal{P}_a = potential of external frcs, V = multi-body interaction potential

$$\mathcal{E}^a(y) = \sum_{x \in \mathcal{L}} V(y(x+r) - y(x); r \in \mathcal{R})$$

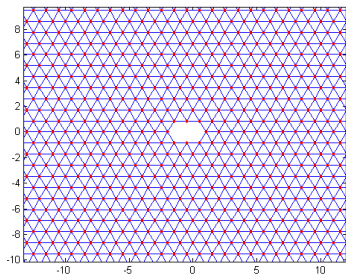
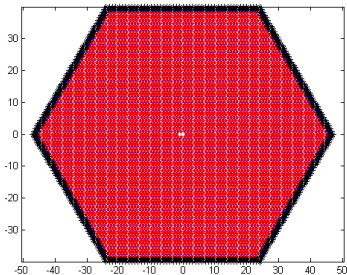
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Example: microcrack

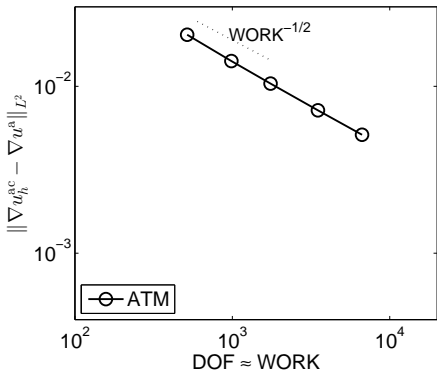
Atomistic Simulation

microcrack under macroscopic shear/stretch, EAM potential

$$V(g) := \sum_{\rho \in \mathcal{R}} \phi(|g_\rho|) + G \left(\sum_{\rho \in \mathcal{R}} \psi(|g_\rho|) \right), \quad \text{where}$$

$$\phi(s) := e^{-2A(s-1)} - 2e^{-A(s-1)}, \quad \psi(s) := e^{-Bs},$$

$$\text{and } G(s) := C((s - s_0)^2 + (s - s_0)^4).$$



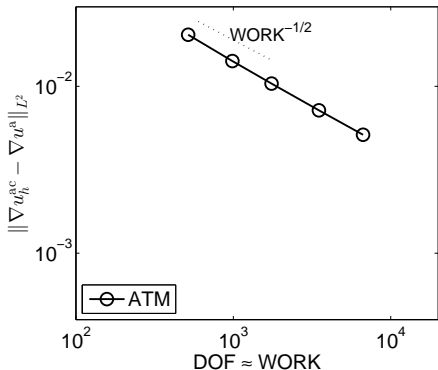
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Can we find a better method,
e.g., accuracy scales $O(N^{-1})$?

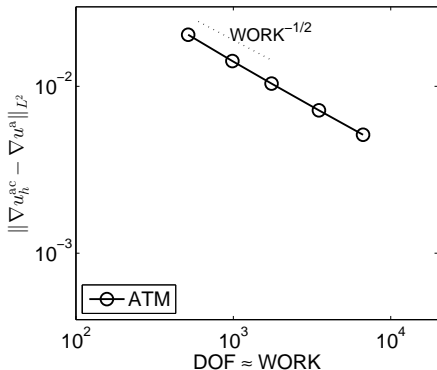
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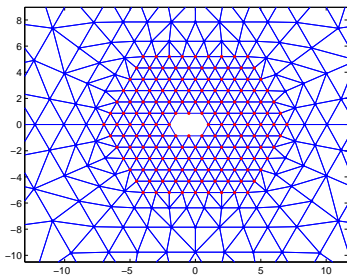
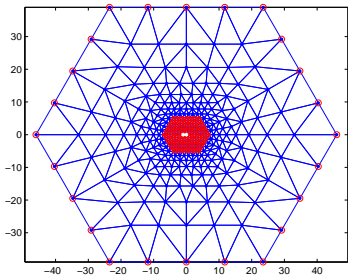


Can we find a better method, e.g., accuracy scales $O(N^{-1})$?

Idea

- Coarse graining the DoF.

Coarse Graining: Adaptive FEM



- $\|\nabla y_a - \nabla y_h\| \leq \|h\nabla^2 y_a\|_{\Omega \setminus \Omega_a}$ [Lin, Ortner/Süli, Lin/Shapeev, ...]

But complexity to evaluate $\mathcal{E}^a|_{\mathcal{Y}_h}$, $\delta\mathcal{E}^a|_{\mathcal{Y}_h}$ is still $O(N)$!

Coarse Graining: Cauchy–Born Approximation

Atomistic Stored Energy:

$$\mathcal{E}^a(y) = \sum_{x \in \mathcal{L}} V(y(x+r) - y(x); r \in \mathcal{R})$$

Cauchy–Born Stored Energy:

$$\mathcal{E}^c(y) = \int_{\Omega} W(\nabla y) \, dV, \quad \text{where } W(F) = V(\{Fr; r \in \mathcal{R}\}).$$

Theorem:

[2007, E/Ming]

Let $y_a \in \operatorname{argmin} \mathcal{E}_a^{\text{tot}}$ be “sufficiently smooth globally”, then there exists $y_c \in \operatorname{argmin} \mathcal{E}_c^{\text{tot}}$ such that

$$\|\nabla y_a - \nabla y_c\|_{L^2} \lesssim C(\|\nabla^3 y_a\|_{L^2} + \|\nabla^2 y_a\|_{L^4}^2)$$

Coarse Graining: Cauchy–Born Approximation

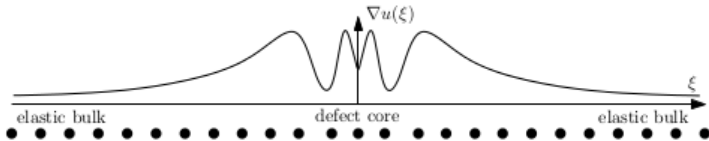
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Cauchy–Born Stored Energy:

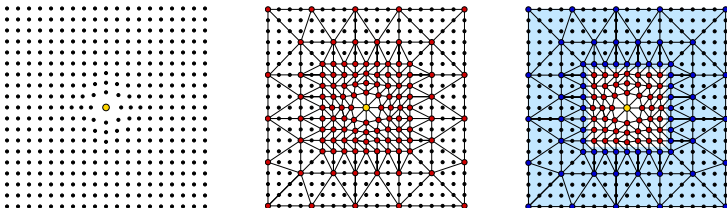
$$\mathcal{E}^c(y) = \int_{\Omega} W(\nabla y) dV, \quad \text{where } W(F) = V(\{Fr; r \in \mathcal{R}\}).$$

- If there are no defects, then the Cauchy–Born model is a highly accurate continuum approximation, $\|y_a - y_c\|_{H^1} \sim N^{-2}$.
- If there are defects, then the Cauchy–Born model has $O(1)$ error.



[2012, Ortner & Luskin]

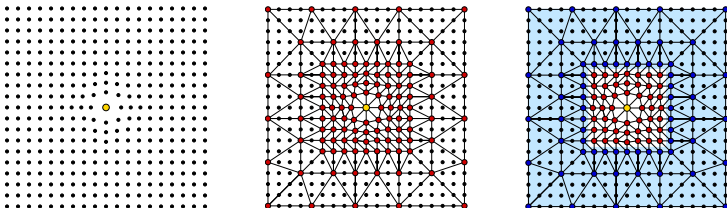
Atomistic/Continuum Coupling: First Attempt



$$\mathcal{E}^a(y_h) \approx \mathcal{E}^{\text{qce}}(y_h) := \sum_{x \in \mathcal{L}_a} \omega_x V_x + \int_{\Omega_c} W(\nabla y_h) dx$$

[Tadmor, Ortiz, Philips 1996]

Atomistic/Continuum Coupling: First Attempt



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[Tadmor, Ortiz, Philips 1996]

Fails the **patch test (ghost force)**:

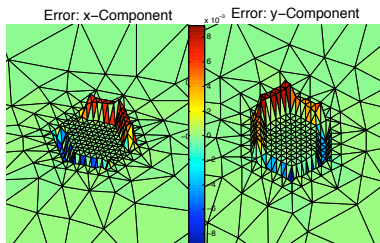
$$\delta \mathcal{E}^a(y_F) = 0$$

and

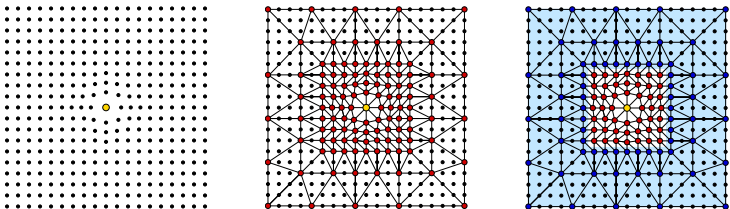
$$\delta \mathcal{E}^c(y_F) = 0,$$

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$$\delta \mathcal{E}^{\text{qce}}(y_F) \neq 0 !$$



Atomistic/Continuum Coupling: First Attempt



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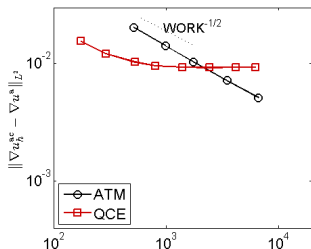
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and

$$\delta \mathcal{E}^c(y_F) = 0,$$

but

$$\delta \mathcal{E}^{\text{qce}}(y_F) \neq 0 !$$



A 1D Model Problem

- Periodic displacements:

$$\mathcal{U} = \{\mathbf{u} = (u_n)_{n \in \mathbb{Z}} : u_{n+N} = u_n, \sum_{n=1}^N u_n = 0\},$$

$$\mathcal{Y} = \{\mathbf{y} = (y_n)_{n \in \mathbb{Z}} : y_n = x_n + u_n \text{ where } \mathbf{u} \in \mathcal{U}\}.$$

- Atomistic energy: (next-nearest neighbor pair interactions)

$$\mathcal{E}^a(\mathbf{y}) = \sum_{n=1}^N \phi(y'_n) + \sum_{n=1}^N \phi(y'_n + y'_{n+1}) = \sum_{n=1}^N \mathcal{E}_n^a(\mathbf{y})$$

where $\mathcal{E}_n^a(\mathbf{y}) = \frac{1}{2} \{\phi(y'_{n-1} + y'_n) + \phi(y'_n) + \phi(y'_{n+1}) + \phi(y'_n + y'_{n+1})\}$

- Continuum finite element model

$$\mathcal{E}^c(\mathbf{y}) = \sum_{n=1}^N \{\phi(y'_n) + \phi(2y'_n)\} = \sum_{n=1}^N \mathcal{E}_n^c(\mathbf{y})$$

where $\mathcal{E}_n^c(\mathbf{y}) = \frac{1}{2} \{\phi(2y'_n) + \phi(y'_n) + \phi(y'_{n+1}) + \phi(2y'_{n+1})\}$

The Energy-Based Quasicontinuum Method

- Choose atomistic and continuum regions:

$$\mathcal{N}^a \cup \mathcal{N}^c = 1, \dots, N$$

- Define a/c hybrid energy

$$\mathcal{E}^{\text{qce}}(y) = \sum_{n \in \mathcal{N}^a} \mathcal{E}_n^a(y) + \underbrace{\sum_{n \in \mathcal{N}^c} \mathcal{E}_n^c(y)}_{\int_{\Omega^c} W(Dy) dx} - \langle g, y \rangle$$



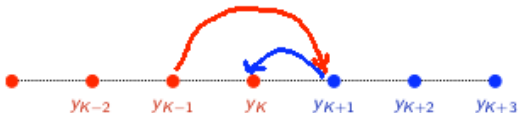
Ghost Forces

Solutions for \mathcal{E}^a and \mathcal{E}^c :

$$\nabla \mathcal{E}^a(x) = 0 \quad \text{and} \quad \nabla \mathcal{E}^c(x) = 0$$

Insert $y_a = x$ into $\nabla \mathcal{E}^{\text{qce}}$

$$\frac{\partial \mathcal{E}^{\text{qce}}}{\partial y_n} \Big|_{y=x} = \frac{\phi'(2)}{2} \times \begin{cases} 0, & n = \dots, K-2 \\ 1, & n = K-1 \\ -1, & n = K \\ -1, & n = K+1 \\ 1, & n = K+2 \\ 0, & n = K+3, \dots \end{cases}$$



Alternative Approaches

1 Energy-based coupling: Interface Correction

2 Force-based coupling:

- FeAt: Kohlhoff, Schmauder, Gumbsch (1989, 1991)
- Dead-load GF removal: Shenoy, Miller, Rodney, Tadmor, Phillips, Ortiz (1999)
- CADD: Shilkrot, Curtin, Miller (2002, ...)
- ...

3 Blending methods: $E = \chi E_a + (1 - \chi) E_c$

- Belytschko & Xiao (2004)
- Parks, Gunzburger, Fish, Badia, Bochev, Lehoucq, et al. (2008)
- ...

4 Quadrature approaches

- Knapp, Ortiz, Gunzburger, ...

Consistent Energy-Based Coupling

Problem: "... the disadvantage of the energy based approach is that it is **extremely difficult** to eliminate the non-physical side effects of the coupled energy functional, dubbed 'ghost forces'."

– Tadmor & Miller, 2009

Goal: Construct consistent A/C energy \mathcal{E}^{ac} by interface correction

$$\mathcal{E}^{\text{ac}}(y) = \sum_{x \in \mathcal{L}_a} V_x + \sum_{x \in \mathcal{L}_i} \tilde{V}_x + \int_{\Omega_c} W(Dy) dx$$

find \tilde{V} s.t. patch test consistency holds: $\delta \mathcal{E}^{\text{ac}}(y_F) = 0$ for all $F \in \mathbb{R}^{d \times d}$.

Questions:

- 1 Does patch test consistency implies accuracy? A priori analysis?
- 2 How to construct consistent coupling method?

A Priori Error Analysis

Framework: Let $y_a \in \operatorname{argmin} \mathcal{E}_a^{\text{tot}}$, $y_{ac} \in \operatorname{argmin} \mathcal{E}_{ac}^{\text{tot}}$, then

$$\|\nabla(y_a - y_{ac})\|_{L^2} \approx \frac{\text{CONSISTENCY}}{\text{STABILITY}} = \frac{\|\delta\mathcal{E}^a(y_a) - \delta\mathcal{E}^{ac}(y_a)\|_{H^{-1}}}{\inf_{\|\nabla u\|_{L^2}=1} \langle \delta^2\mathcal{E}^{ac}(y_a)u, u \rangle}$$

3 Steps:

- 1 **CONSISTENCY:** $\langle \delta\mathcal{E}^{ac}(y_a) - \delta\mathcal{E}^a(y_a), u_h \rangle \lesssim h \|\nabla^2 y_a\|_{L^2(\Omega_c)} \|\nabla u_h\|_{L^2}$
- 2 **STABILITY:** $\langle \delta^2\mathcal{E}^{ac}(y)u, u \rangle \geq C_{\text{stab}} \|\nabla u\|_{L^2}^2$
- 3 **REGULARITY:** bounds on $\nabla^2 y_a$, note that r^{-a} for defects ($a = 2$ dislocations, $a = 3$ vacancy)

Numerical Analysis Literature on A/C Coupling

cf 'Atomistic-to-Continuum Coupling', Acta Numerica (2013), Luskin/Ortner

Analysis for Energy-Based Coupling:

- 1D
 - NN, variational analysis: [Blanc/LeBris/Legoll, 2005]
 - Error estimates for QCE and QNL: [Dobson/Luskin, 2009], [Ming/Yang, 2010], [Ortner, 2010], [Ortner/Wang, 2011], [Li/Luskin, 2011]
 - Sharp Stability Analysis, Linear Regime: [Dobson/Luskin/Ortner, 2010]
 - blending methods: [Luskin/van Koten, 2012]
- 2D, pair interactions, defects: [Ortner/Shapeev, 2012]
- 2D, first order consistency for general finite-range interactions: [Ortner, 2012]
- 3D, pair interactions: [Shapeev, 2012]
- 2D, multi-body interactions: [Ortner/LZ, 2012, 2013]
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- 1D/2D, stability of consistent energy based coupling: [Ortner/Shapeev/LZ, 2013]

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Analysis for Multi-Lattices:

Dobson/Elliott/Luskin/Tadmor (2007); Abdulle/Lin/Shapeev (2011, 2012)

Patch Test Consistency implies First-order Consistency?

Suppose \mathcal{E}^{ac} is patch test consistent (no ghost force for homogeneous deformation):

$$\delta\mathcal{E}^{\text{ac}}(y_F) = 0 \quad \forall F \in \mathbb{R}^{d \times d}$$

“Theorem:”

[First-order Consistency]

Suppose $\delta\mathcal{E}^{\text{ac}}$ passes the patch test, V finite range multi-body potential + technical conditions +

- $d = 1$; or
- $d = 2$, Ω_a connected; [Ortner, 2012] or
- $d = 3$, Ω_a connected [LZ, 2013]

then

$$\langle \delta\mathcal{E}^{\text{ac}}(y) - \delta\mathcal{E}^{\text{a}}(y), u_h \rangle \lesssim \|h\nabla^2 y\|_{L^2(\Omega_c \cup \Omega_i)} \|\nabla u_h\|_{L^2}$$

With the assumption of stability, $\|y_a - y_{\text{ac}}\| \sim N^{-1}$.

Consequence of Patch Test Consistency

If an A/C energy \mathcal{E}^{ac} satisfies patch test consistency,

$$0 = \langle \delta \mathcal{E}^{\text{ac}}(y_{\text{F}}), u \rangle = \sum_{T \in \mathcal{T}} |T| \Sigma_{\text{ac}}(y_{\text{F}}; T) : \nabla_T u$$

then Σ_{ac} is **discrete divergence free**.

Lemma:

\exists a function $\psi(\mathbf{F}, T) \in N_1(\mathcal{T})^2$, such that

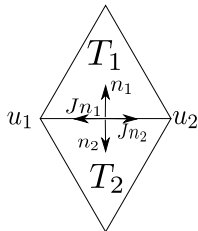
$$\Sigma_{\text{ac}}(y_{\text{F}}; T) = \partial W(\mathbf{F}) + J \nabla \psi(\mathbf{F}; T)$$

$N_1(\mathcal{T})$ is **Crouzeix–Raviart** finite element space,

J is the counter-clockwise rotation by $\pi/2$.

$J \nabla \psi(\mathbf{F}; T)$ is divergence free piecewise constant tensor field

[Arnold/Falk, Polthier/Preuß].



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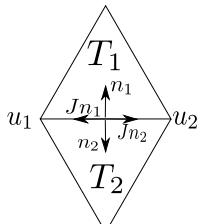
- For general deformation y , deformation gradient average for patch $\omega_f = (T_1 \cup T_2)$, $F_f(y) = \int_{\omega_f} \nabla y \, dx$,

Corrector function: $\hat{\psi}(y; \cdot) = \sum_{f \in \mathcal{F}} \psi(F_f(y); m_f) \zeta_f$

- Define the '**modified**' stress function,

$$\hat{\Sigma}_{\text{ac}}(y; T) := \Sigma_{\text{ac}}(y; T) - J \nabla \hat{\psi}(y; T), \quad \text{for } T \in \mathcal{T}.$$

- $\hat{\Sigma}_{\text{ac}}(y_F; T) = \partial W(F) = \Sigma_a(y_F; T)$



Construction of Consistent A/C Schemes

$$\mathcal{E}^{\text{ac}}(y_h) = \sum_{x \in \mathcal{L}_a} V_x + \sum_{x \in \mathcal{L}_i} \tilde{V}_x + \sum_{x \in \mathcal{L}_c} V_x^c$$

Construct \tilde{V} s.t. $\delta \mathcal{E}^{\text{ac}}(y_F) = 0$ for all $F \in \mathbb{R}^{d \times d}$.

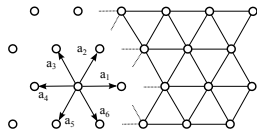
General Construction: [1D, Shimokawa et al, 2004; E/Lu/Yang, 2006]

$$\tilde{V}_x = V(\tilde{g}_{x,r}; r \in \mathcal{R})$$

$$\tilde{g}_{x,r} = \sum_{s \in \mathcal{R}_x} C_{x,r,s} g_s$$

→ Find $C_{x,r,s}$ s.t. $\delta \mathcal{E}^{\text{ac}}(y_F) = 0 \quad \forall F$

→ geometric conditions only!



2d, NN, multibody potential,
triangular lattice

- Explicit constructions for 2D general interface: [Ortner/LZ; SINUM, 2012]
- **In general: compute $C_{x,r,s}$ numerically in preprocessing**

Construction of Consistent A/C Schemes

$$\mathcal{E}^{\text{ac}}(y_h) = \sum_{x \in \mathcal{L}_a} V_x + \sum_{x \in \mathcal{L}_i} \tilde{V}_x + \sum_{x \in \mathcal{L}_c} V_x^c$$

Construct \tilde{V} s.t. $\delta \mathcal{E}^{\text{ac}}(y_F) = 0$ for all $F \in \mathbb{R}^{d \times d}$.

General Construction: [1D, Shimokawa et al, 2004; E/Lu/Yang, 2006]

2. Patch Test Consistency

$$\tilde{V}_x = V(\tilde{g}_{x,r}; r \in \mathcal{R})$$

$$\tilde{g}_{x,r} = \sum_{s \in \mathcal{R}_x} C_{x,r,s} g_s$$

→ Find $C_{x,r,s}$ s.t. $\delta \mathcal{E}^{\text{ac}}(y_F) = 0 \quad \forall F$

→ geometric conditions only!

1. Local Energy Consistency $\tilde{V}(y_F) = V(y_F)$

$$\Rightarrow r = \sum_{s \in \mathcal{R}_x} C_{x,r,s} s. \quad (\text{a})$$

$$0 = \langle \delta \mathcal{E}^{\text{ac}}(y_F), u \rangle$$

$$= \sum_{x \in \mathcal{L}} \sum_{r \in \mathcal{R}} V_{F,r} \sum_{s \in \mathcal{R}} C_{x,r,s} D_s u$$

$$= \sum_{x \in \mathcal{L}} \sum_{r \in \mathcal{R}} \sum_{s \in \mathcal{R}} (C_{x-a_s,r,s} V_{F,r} - C_{x,r,s} D_r V_{F,r}) u(x)$$

$$\Rightarrow \sum_{r \in \mathcal{R}} \sum_{s \in \mathcal{R}} (C_{x-s,r,s} V_{F,r} - C_{x,r,s} V_{F,r}) = 0. \quad (\text{b})$$

Solve (a) + (b) + B.C. in \mathcal{L}_a and \mathcal{L}_c to obtain $C_{x,r,s}$ for $x \in \mathcal{L}_i$.

unknowns: $|\mathcal{L}_I| |\mathcal{R}|^2$, eqns: $|\mathcal{L}_I| (2|\mathcal{R}| + 3)$.

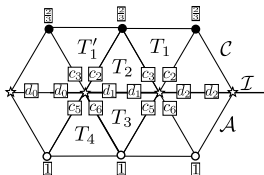
Construction of Consistent A/C Schemes

$$\mathcal{E}^{\text{ac}}(y_h) = \sum_{x \in \mathcal{L}_a} V_x + \sum_{x \in \mathcal{L}_i} \tilde{V}_x + \sum_{x \in \mathcal{L}_c} V_x^c$$

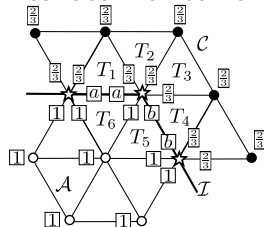
Construct \tilde{V} s.t. $\delta \mathcal{E}^{\text{ac}}(y_F) = 0$ for all $F \in \mathbb{R}^{d \times d}$.

General Construction:

Flat interface



Interface with corner



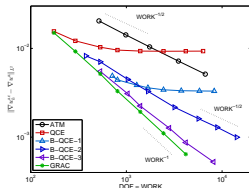
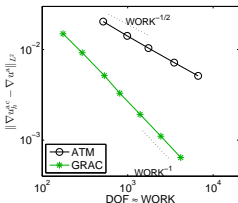
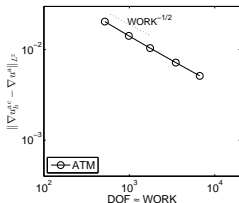
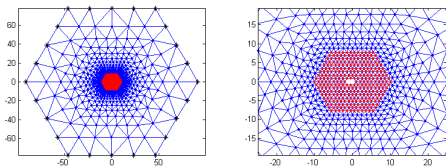
$C_{x,r,r}$ for NN interaction, many-body potential, one-sided construction.

1. works for general interface in 2d
2. preprocessing for longer interaction range

Numerical Experiment

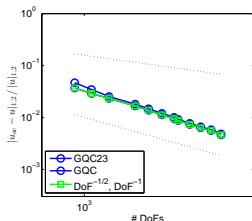
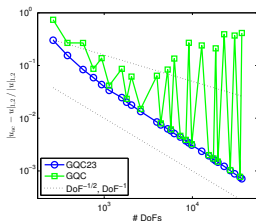
Test Problem: microcrack in the triangular lattice, EAM multi-body potential

$$V = F_\alpha \left(\sum_{i \neq j} \rho_\beta(r_{ij}) \right) + \frac{1}{2} \sum_{i \neq j} \phi_{\alpha\beta}(r_{ij})$$



Optimization of the Reconstruction Coefficients

[Ortner/LZ, 2013]



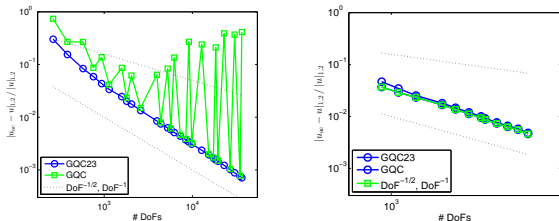
Left: H^1 Error with coefficients from least square solution

Right: H^1 Error with coefficients from L^1 minimization

- The coefficients needs to be pre-computed for longer range interactions, also needs to be optimized for better accuracy.

Optimization of the Reconstruction Coefficients

[Ortner/LZ, 2013]



Left: H^1 Error with coefficients from least square solution

Right: H^1 Error with coefficients from L^1 minimization

- Consistency Error Estimate

$$\begin{aligned} \langle \delta \mathcal{E}^{\text{ac}}(y) - \delta \mathcal{E}^{\text{a}}(y), u_h \rangle &= \sum_{T \in \mathcal{T}} (\Sigma_{\text{ac}}(y; T) - \Sigma_{\text{a}}(y; T)) : \nabla u_h \\ &\leq C \|h \nabla^2 y\|_{L^2(\Omega_c \cup \Omega_i)} \|\nabla u_h\|_{L^2} \end{aligned}$$

The constant C is controlled by $\sum_{r \in \mathbb{R}} \sum_{s \in \mathbb{R}} |r||s| C_{x,r,s}$.

- The coefficients can be obtained by solving a constrained L^1 minimization problem.

Stability of Consistent A/C Coupling Method

[Ortner/Shapeev/LZ, 2013]

- Study the Hessians

$$\langle H_{Dy}^a v, v \rangle := \langle \delta^2 \mathcal{E}^a(y) v, v \rangle := \sum_{\xi \in \mathbb{Z}} \sum_{\xi, \varsigma \in \mathcal{R}} V_{\rho\varsigma}(Dy(\xi)) \cdot D_\rho v(\xi) D_\varsigma v(\xi)$$

$$\langle H_{Dy}^{ac} v, v \rangle := \langle \delta^2 \mathcal{E}^{ac}(y) v, v \rangle := \sum_{\xi \in \mathbb{Z}} \sum_{\xi, \varsigma \in \mathcal{R}} \tilde{V}_{\rho\varsigma}(Dy(\xi)) \cdot D_\rho v(\xi) D_\varsigma v(\xi)$$

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- Stability constant:

$$\gamma(H) := \inf_{\substack{u \in \mathcal{W}_0 \\ \|\nabla u\|_{L^2} = 1}} \langle Hu, u \rangle.$$

We say that H is stable if $\gamma(H) > 0$.

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- For homogenous deformation y_F ,
 - $\gamma(H_F^{ac}) \leq \gamma(H_F^a)$ for all $F > 0$.
 - $\gamma(H_F^c) = W''(F) \geq \gamma(H_F^a)$ for all $F > 0$.

Universally Stable Method

Question: For any potential V , can we find such a A/C scheme, such that $\gamma_F^{\text{ac}} > 0$ if and only if $\gamma_F^{\text{a}} > 0$? If exists, such method is called universally stable.

- universally stable method in 1D

$$z^* := \begin{cases} z(\xi), & \xi \leq 0, \\ 2z(0) - z(-\xi), & \xi > 0. \end{cases}$$

$$\mathcal{E}^{\text{rfl}}(y) := \mathcal{E}^*(y) + \int_0^\infty W(\nabla y) \, dx, \quad \text{where}$$

$$\mathcal{E}^*(y) := \sum_{\xi=-\infty}^{-1} [V(Dy^*(\xi)) - V(F\mathcal{R})] + \frac{1}{2} [V(Dy^*(0)) - V(F\mathcal{R})].$$

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- Nonexistence of universally stable method in 2D, even for flat interface.

Stability Gap and Stabilization

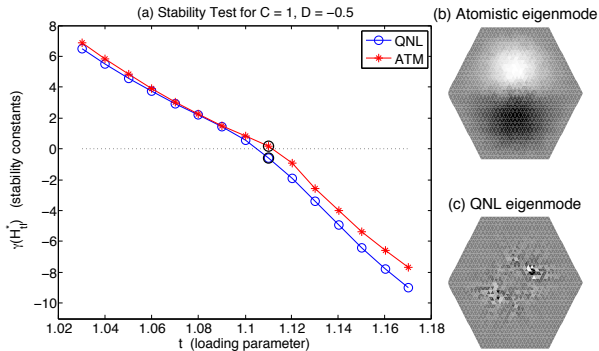


Figure : Stability test for $C = 1, D = -0.5$. The black circles indicate which eigenmodes (u_1 -component) are plotted in (b, c).

$$V(g) := \sum_{\rho \in \mathcal{R}} \phi(|g_\rho|) + G\left(\sum_{\rho \in \mathcal{R}} \psi(|g_\rho|)\right) + D \sum_{j=1}^6 (r_j \cdot r_{j+1} - 1/2)^2,$$

where $\phi(s) := e^{-2A(s-1)} - 2e^{-A(s-1)}$, $\psi(s) := e^{-Bs}$, and $G(s) := C((s - s_0)^2 + (s - s_0)^4)$

Stability Gap and Stabilization

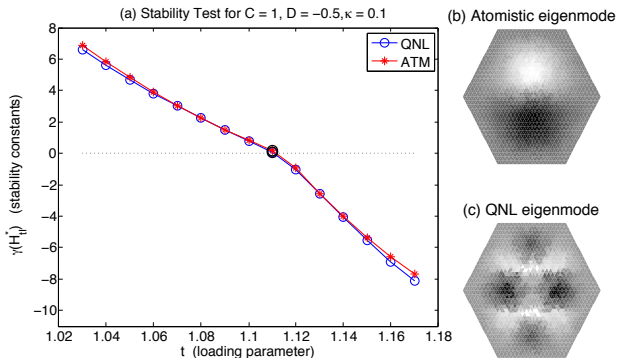


Figure : Stability test for $C = 1, D = -0.5, \kappa = 0.1$. The black circles indicate which eigenmodes (u_1 -component) are plotted in (b, c).

$$\mathcal{E}^{\text{stab}}(y) := \mathcal{E}^{\text{ac}}(y) + \kappa \langle Su, u \rangle, \quad \text{for } y = Fx + u, u \in \mathcal{W}_0,$$

where

$$\langle Su, u \rangle := \sum_{\xi \in \mathcal{L}^{(0)}} |D^2 u(\xi)|^2,$$

Stability Gap and Stabilization

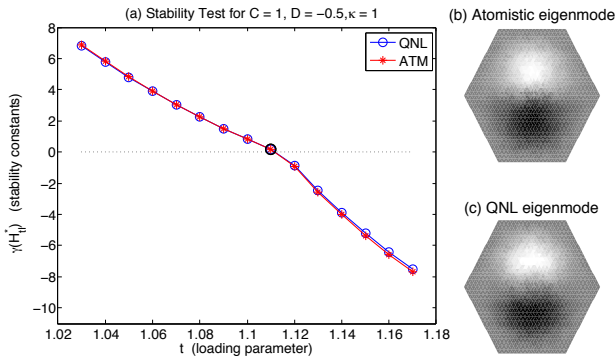


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Stablization: Consistency vs. Stability

“Theorem:”

[Critical Strain for Stabilized A/C Coupling]

Let V have hexagonal symmetry, $F \propto I$, $V_{i,i+2} = V_{i,i+3} \equiv 0$, and $\tilde{c}_1^{(1)} - \tilde{c}_1^{(-1)} \neq 0$; then there exists constants $c_1, c_2 > 0$ such that

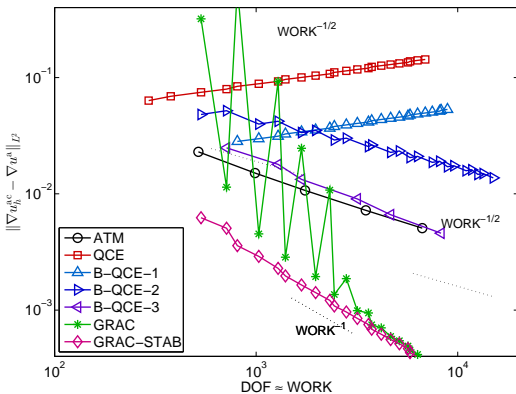
$$\gamma(H_F^a) - \frac{c_1}{\kappa^2} \leq \gamma(H_F^{ac} + \kappa S) \leq \gamma(H_F^a) - \frac{c_2}{\kappa^2}.$$

- existence of a critical loading parameter $t_*^\kappa \in [t_0, t_*]$ for which $\gamma(H_{t_*^\kappa}^{ac} + \kappa S) = 0$ and such that

$$|t_*^\kappa - t_*| \approx \frac{1}{\kappa^2}.$$

- Therefore, if we wish to admit at most an $O(N^{-1})$ error in the critical strain, then we must accordingly choose $\kappa = O(N^{1/2})$. Unfortunately, this has a consequence for the consistency error of the stabilised A/C method, which will accordingly scale like $O(N^{1/2})$.

Recipe for General Interaction Range



GRAC-STAB =
 Geometric Reconstruction +
 L^1 Minimization +
 Stabilization

2nd nearest neighbor EAM potential

Outlook

Summary

- Consistent energy based method has accuracy $O(N^{-1})$
- Construction of consistent energy-based a/c methods
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- Optimize the A/C coupling method
- Implementation, benchmarks, applications
- 3D
- dislocation

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- A/C methods for multi-lattices
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- A/C methods for molecular dynamics

Thanks for Your Attention!

The Force-based Approach

A possible solution are **force-based a/c methods**:

- Dead-load GF removal: Shenoy, Miller, Rodney, Tadmor, Phillips, Ortiz (1999)
- FeAt: Kohlhoff, Schmauder, Gumbsch (1989, 1991)
- AtC: Parks, Gunzburger, Fish, Badia, Bochev, Lehoucq, et al. (2007, ...)
- CADD: Shilkrot & Curtin & Miller (2002, ...)
- Adaptive Resolution MD: Delle Site et al. (2007, ...)
- Brutal Force Mixing: Bernstein, Csanyi et al. (2007, ...)
-
- Analysis: Dobson & Luskin (2008); Ming (2009); Dobson & Luskin & Ortner (2009, 2010, 2010); Makridakis & Ortner & Süli (2010, preprint); Dobson & Ortner & Shapeev (preprint), Lu & Ming (manuscript)

Consistency

Numerical Analysis Literature on Energy-Based Coupling:

- 1D, NN, variational analysis: [Blanc/LeBris/Legoll, 2005]
- 1D, Error estimates for QCE and QNL: [Dobson/Luskin, 2009], [Ming/Yang, 2010], [Ortner, 2010], [Ortner/Wang, 2011]
- 1D, Sharp Stability Analysis, Linear Regime: [Dobson/Luskin/Ortner, 2010]
- 1D, finite range: [Li/Luskin, preprint]
- 1D, EAM potentials: [Li/Luskin, 2011]
- 1D, blending methods: [Luskin/van Koten, preprint]
- 2D, pair interactions, defects: [Ortner/Shapeev, preprint]
- 2D, NN interactions: [Ortner/Zhang, manuscript]
- 2D, consistency for general finite-range interactions: [Ortner, preprint]

Cauchy–Born Approximation

Atomistic Stored Energy:

$$\mathcal{E}^a(y) = \sum_{x \in \mathcal{L}} V(y(x+r) - y(x); r \in \mathcal{R})$$

Cauchy–Born Stored Energy:

$$\mathcal{E}^c(y) = \int_{\Omega} W(\nabla y) \, dV, \quad \text{where } W(F) = V(\{Fr; r \in \mathcal{R}\}).$$

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Theorem:

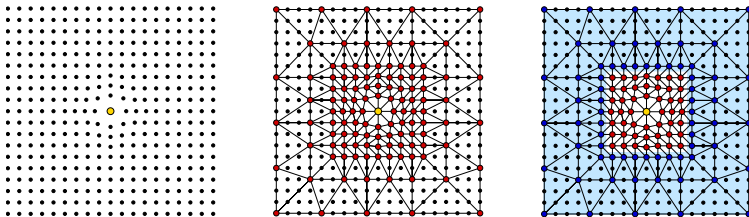
[Similar to result by E/Ming; 2007]

Let $y_a \in \operatorname{argmin} \mathcal{E}_a^{\text{tot}}$ be “sufficiently smooth globally”, then there exists $y_c \in \operatorname{argmin} \mathcal{E}_c^{\text{tot}}$ such that

$$\|\nabla y_a - \nabla y_c\|_{L^2} \lesssim C(\|\nabla^3 y_a\|_{L^2} + \|\nabla^2 y_a\|_{L^4}^2)$$

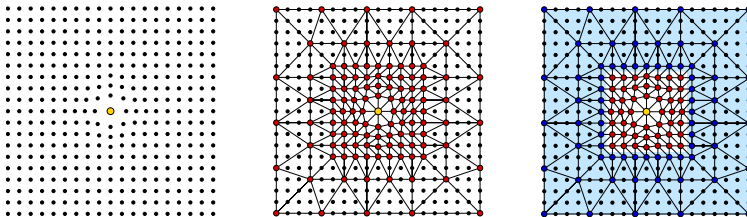
If there are no defects, then the Cauchy–Born model is a highly accurate continuum approximation.

Atomistic/Continuum Coupling: First Attempt



$$\mathcal{E}^a(y_h) \approx \mathcal{E}^{ac}(y_h) := \sum_{x \in \mathcal{L}_a} V_x + \int_{\Omega_c} W(\nabla y_h) dx$$

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Fails the **patch test**:

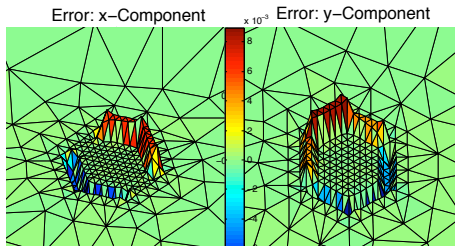
$$\delta \mathcal{E}^a(y_F) = 0$$

and

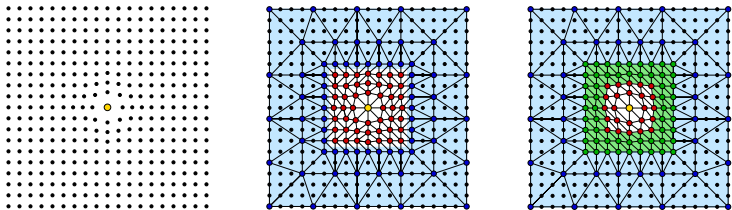
$$\delta \mathcal{E}^c(y_F) = 0,$$

but

$$\delta \mathcal{E}^{ac}(y_F) \neq 0$$



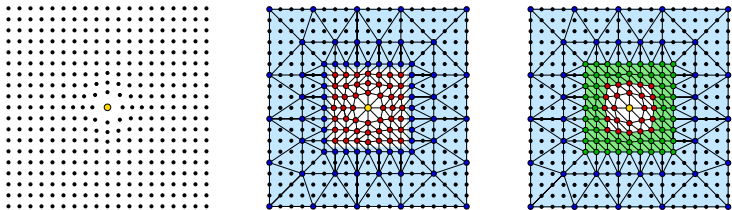
A/C Coupling: Interface Correction



$$\mathcal{E}^{\text{ac}}(y_h) = \sum_{x \in \mathcal{L}_a} V_x + \sum_{x \in \mathcal{L}_i} \tilde{V}_x + \int_{\Omega_c} W(\nabla y_h) dx$$

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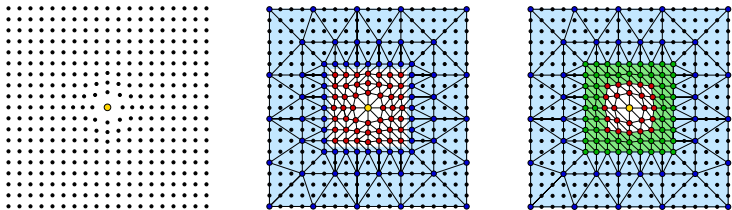
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Main Challenges:

- **Construction** of patch test consistent \mathcal{E}^{ac} : still largely unsolved

[Shimokawa et al; 2004], [E/Lu/Yang; 2006], [Shapeev; preprint], [Iyer/Gavini; preprint], [Ortner/Zhang; preprint], [Xiao/Belytschko, 2004], [Klein/Zimmermann, 2006], ...

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- **Accuracy**: Does passing the patch test imply “good” accuracy?

Consistency of The Schemes

$$\mathcal{E}^{\text{ac}}(y_h) = \sum_{x \in \mathcal{L}_a} V_x + \sum_{x \in \mathcal{L}_i} \tilde{V}_x + \sum_{x \in \mathcal{L}_c} V_x^c$$

Construct \tilde{V} s.t. $\delta \mathcal{E}^{\text{ac}}(y_F) = 0$ for all $F \in \mathbb{R}^{d \times d}$.

$$\tilde{V}_x = V(\tilde{g}_{x,r}; r \in \mathcal{R})$$

$$\tilde{g}_{x,r} = \sum_{s \in \mathcal{R}_x} C_{x,r,s} g_s$$

→ Find $C_{x,r,s}$ s.t. $\delta \mathcal{E}^{\text{ac}}(y_F) = 0 \quad \forall F$

→ geometric conditions only!

Theorem: [Ortner/Zhang, '11]

There exists a constant C that depends only on $M_2(V)$ and $M_3(V)$ such that

$$\|\delta \mathcal{E}^{\text{ac}}(y) - \delta \mathcal{E}^{\text{a}}(y)\|_{\mathcal{W}^{-1,p}} \leq C(\|D^3 y\|_{\ell^p(\Omega_c)} + \|D^2 y\|_{\ell^{2p}(\Omega_c)}^2 + \|D^2 y\|_{\ell^p(\Omega_I)})$$

Outlook on A/C Methods

Summary

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- A/C methods for Coulomb interaction
- A/C methods for electronic structure models
(done only for insulators)
- A/C methods for molecular dynamics

Consequence of Patch Test Consistency

If an A/C energy \mathcal{E}^{ac} satisfies patch test consistency,

$$0 = \langle \delta \mathcal{E}^{\text{ac}}(y_F), u \rangle = \sum_{T \in \mathcal{T}} |T| \Sigma_{\text{ac}}(y_F; T) : \nabla_T u$$

then Σ_{ac} is **discrete divergence free**.

Lemma:

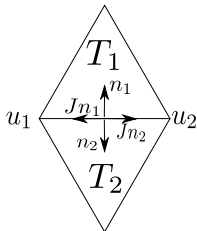
\exists a function $\psi(F, T) \in N_1(\mathcal{T})^2$, such that

$$\Sigma_{\text{ac}}(y_F; T) = \partial W(F) + J \nabla \psi(F; T)$$

$N_1(\mathcal{T})$ is **Crouzeix–Raviart** finite element space,

J is the counter-clockwise rotation by $\pi/2$.

$J \nabla \psi(F; T)$ is divergence free piecewise constant tensor field
[Arnold/Falk, Polthier/Preuß].



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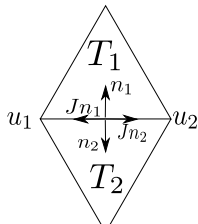
- For general deformation y , deformation gradient average for patch $\omega_f = (T_1 \cup T_2)$, $F_f(y) = \int_{\omega_f} \nabla y \, dx$,

Corrector function: $\hat{\psi}(y; \cdot) = \sum_{f \in \mathcal{F}} \psi(F_f(y); m_f) \zeta_f$

- Define the 'modified' stress function,

$$\hat{\Sigma}_{\text{ac}}(y; T) := \Sigma_{\text{ac}}(y; T) - J \nabla \hat{\psi}(y; T), \quad \text{for } T \in \mathcal{T}.$$

- $\hat{\Sigma}_{\text{ac}}(y_F; T) = \partial W(F) = \Sigma_a(y_F; T)$



Consistency of The Schemes

$$\mathcal{E}^{\text{ac}}(y_h) = \sum_{x \in \mathcal{L}_a} V_x + \sum_{x \in \mathcal{L}_i} \tilde{V}_x + \sum_{x \in \mathcal{L}_c} V_x^c$$

Construct \tilde{V} s.t. $\delta \mathcal{E}^{\text{ac}}(y_F) = 0$ for all $F \in \mathbb{R}^{d \times d}$.

“Theorem:”

[Ortner/Zhang]

There exists a constant C that depends only on $M_2(V)$ and $M_3(V)$ such that

$$\|\delta \mathcal{E}^{\text{ac}}(y) - \delta \mathcal{E}^{\text{a}}(y)\|_{\mathcal{W}^{-1,p}} \leq C(\|\nabla^3 y\|_{\ell^p(\Omega_c)} + \|\nabla^2 y\|_{\ell^{2p}(\Omega_c)}^2 + \|\nabla^2 y\|_{\ell^p(\Omega_I)}).$$

Proof of Consistency

$$\begin{aligned}\langle \delta \mathcal{E}^{\text{ac}}(y) - \delta \mathcal{E}^{\text{a}}(y), u \rangle &= \sum_{T \in \mathcal{T}} (\Sigma_{\text{ac}}(y; T) - \Sigma_{\text{a}}(y; T)) : \nabla u \\ &= \sum_{T \in \mathcal{T}} (\widehat{\Sigma}_{\text{ac}}(y; T) - \Sigma_{\text{a}}(y; T)) : \nabla u\end{aligned}$$

- 1 $T \in \Omega_{\mathcal{A}}, \widehat{\Sigma}_{\text{ac}}(y; T) = \Sigma_{\text{a}}(y; T),$
- 2 $T \in \Omega_{\mathcal{C}}, \widehat{\Sigma}_{\text{ac}}(y; T) - \Sigma_{\text{a}}(y; T) = \Sigma_{\mathcal{C}}(y; T) - \Sigma_{\text{a}}(y; T),$ 2nd order consistency
- 3 $T \in \Omega_{\mathcal{I}},$ Let $y_T = y_{\nabla y(T)},$ we have,

$$\begin{aligned}|\widehat{\Sigma}_{\text{ac}}(y; T) - \Sigma_{\text{a}}(y; T)| &\leq |\widehat{\Sigma}_{\text{ac}}(y; T) - \partial W(\nabla y(T))| + |\partial W(\nabla y(T)) - \Sigma_{\text{a}}(y; T)| \\ &= |\widehat{\Sigma}_{\text{ac}}(y; T) - \widehat{\Sigma}_{\text{ac}}(y_T; T)| + |\Sigma_{\text{a}}(y_T; T) - \Sigma_{\text{a}}(y; T)| \\ &\leq C \|\nabla y(T) - \nabla y\|_{\ell^\infty} \leq C |D^2 y|\end{aligned}$$

\Rightarrow

$$\|\delta \mathcal{E}^{\text{ac}}(y) - \delta \mathcal{E}^{\text{a}}(y)\|_{\mathcal{W}^{-1,p}} \leq C (\|\nabla^3 y\|_{\ell^p(\Omega_{\mathcal{C}})} + \|\nabla^2 y\|_{\ell^{2p}(\Omega_{\mathcal{C}})}^2 + \|\nabla^2 y\|_{\ell^p(\Omega_{\mathcal{I}})}).$$